## Finite Element Modeling

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#### Two/Three-dimensional Finite Element Method

## Finite Element Method (FEM)

#### inflatable link simulation



S, Mises SNEG, (fraction = -1.0) (Avg: 75%)



### Finite Element Method (FEM) Integral forms (weak forms) apply

strain potential energy (one-dim. beam)

$$U = \int_0^L \frac{1}{2} EA\left(\frac{\partial u}{\partial x}\right)^2 \mathrm{d}x$$

kinetic energy (one-dim. beam)

$$T = \int_0^L \frac{1}{2} \rho A \left(\frac{\partial u}{\partial t}\right)^2 \mathrm{d}x$$

How calculate energies in integral forms?

# Finite Element Method (FEM)

$$\int_0^L = \int_{x_1}^{x_2} + \int_{x_2}^{x_3} + \int_{x_3}^{x_4} + \int_{x_4}^{x_5}$$

apply piecewise linear approximation

$$\int_{x_i}^{x_j} = \frac{1}{2} \begin{bmatrix} u_i & u_j \end{bmatrix} \begin{bmatrix} & & \\ & & \end{bmatrix} \begin{bmatrix} u_i \\ u_j \end{bmatrix}$$

synthsize

$$\int_0^L = \frac{1}{2} \begin{bmatrix} u_1 & u_2 & \cdots & u_5 \end{bmatrix} \begin{bmatrix} u_1 & u_2 & \cdots & u_5 \end{bmatrix}$$

one-dimensional deformation extensional strain  $\varepsilon$ Young's modulus Estrain potential energy density  $\frac{1}{2}E\varepsilon^2$ 

extensional & shear strains  $\rightarrow$  strain vector  $\varepsilon$ Lamé's constants  $\lambda$ ,  $\mu \rightarrow$  elasticity matrix  $\lambda I_{\lambda} + \mu I_{\mu}$ strain potential energy density  $\frac{1}{2}\varepsilon^{T}(\lambda I_{\lambda} + \mu I_{\mu})\varepsilon$ 



natural state moved and deformed state

displacement vector

$$\boldsymbol{u}(x,y) = \left[\begin{array}{c} \boldsymbol{u}(x,y)\\ \boldsymbol{v}(x,y) \end{array}\right]$$



natural deformed and rotated



shear deformation

rotational motion

$$\frac{\partial u}{\partial x} = \text{extension along } x \text{-axis} \quad \frac{\partial v}{\partial y} = \text{extension along } y \text{-axis}$$
$$\frac{\partial v}{\partial x} = \text{shear} + \text{rotation} \qquad \frac{\partial u}{\partial y} = \text{shear} - \text{rotation}$$

Cauchy strain

$$\varepsilon_{xx} = \frac{\partial u}{\partial x}, \qquad \varepsilon_{yy} = \frac{\partial v}{\partial y}, \qquad 2\varepsilon_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$$

 $\downarrow$ 



Strain potential energy density linear isotropic elastic material

$$rac{1}{2}oldsymbol{arepsilon}^{\mathrm{T}}(\lambda I_{\lambda}+\mu I_{\mu})oldsymbol{arepsilon}$$

where  $\lambda$  and  $\mu$  are Lamé's constants and

$$I_\lambda = \left[egin{array}{cccc} 1 & 1 \ 1 & 1 \ \end{array}
ight], \quad I_\mu = \left[egin{array}{ccccc} 2 & \ & 2 \ & & 1 \end{array}
ight]$$

Volume element

$$h\,\mathrm{d}S = h\,\mathrm{d}x\,\mathrm{d}y$$

## Strain potential energy $U = \int_{S} \frac{1}{2} \varepsilon^{\mathrm{T}} (\lambda I_{\lambda} + \mu I_{\mu}) \varepsilon \ h \, \mathrm{d}S$

Volume element

$$h\,\mathrm{d}S = h\,\mathrm{d}x\,\mathrm{d}y$$

Strain potential energy  
$$U = \int_{S} \frac{1}{2} \varepsilon^{\mathrm{T}} (\lambda I_{\lambda} + \mu I_{\mu}) \varepsilon \ h \, \mathrm{d}S$$

Kinetic energy

$$T = \int_{S} \frac{1}{2} \rho \, \dot{\boldsymbol{u}}^{\mathrm{T}} \dot{\boldsymbol{u}} \, h \, \mathrm{d}S$$



displacement vector

$$\boldsymbol{u}(x,y,z) = \begin{bmatrix} u(x,y,z) \\ v(x,y,z) \\ w(x,y,z) \end{bmatrix}$$



uvw
$$\partial/\partial x$$
ext. along xshr + rot in xyshr - rot in zx $\partial/\partial y$ shr - rot in xyext. along yshr + rot in yz $\partial/\partial z$ shr + rot in zxshr - rot in yzext. along z

$$2 \cdot \text{shear in } yz\text{-plane} = \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y}$$
$$2 \cdot \text{shear in } zx\text{-plane} = \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z}$$
$$2 \cdot \text{shear in } xy\text{-plane} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$$

#### Cauchy strain

$$\varepsilon_{xx} = \frac{\partial u}{\partial x} \quad \varepsilon_{yy} = \frac{\partial v}{\partial y} \quad \varepsilon_{zz} = \frac{\partial w}{\partial z}$$
$$2\varepsilon_{yz} = \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y}$$
$$2\varepsilon_{zx} = \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z}$$
$$2\varepsilon_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$$



Strain potential energy density linear isotropic elastic material

$$rac{1}{2}oldsymbol{arepsilon}^{\mathrm{T}}(\lambda I_{\lambda}+\mu I_{\mu})oldsymbol{arepsilon}$$



#### Volume element

$$\mathrm{d}V = \mathrm{d}x\,\mathrm{d}y\,\mathrm{d}z$$

## Strain potential energy $U = \int_{V} \frac{1}{2} \boldsymbol{\varepsilon}^{\mathrm{T}} (\lambda I_{\lambda} + \mu I_{\mu}) \boldsymbol{\varepsilon} \, \mathrm{d}V$

#### Volume element

$$\mathrm{d}V = \mathrm{d}x\,\mathrm{d}y\,\mathrm{d}z$$

## Strain potential energy $U = \int_{V} \frac{1}{2} \boldsymbol{\varepsilon}^{\mathrm{T}} (\lambda I_{\lambda} + \mu I_{\mu}) \boldsymbol{\varepsilon} \, \mathrm{d}V$

Kinetic energy

$$T = \int_{V} \frac{1}{2} \rho \; \dot{\boldsymbol{u}}^{\mathrm{T}} \dot{\boldsymbol{u}} \; \mathrm{d}V$$





assume density  $\rho$  and thickness h are constants kinetic energy of  $\triangle = \triangle P_i P_j P_k$ 

$$T_{i,j,k} = \int_{\Delta} \frac{1}{2} \rho \, \dot{\boldsymbol{u}}^{\mathrm{T}} \, \dot{\boldsymbol{u}} \, h \, \mathrm{d}S$$
  
=  $\frac{1}{2} \begin{bmatrix} \dot{\boldsymbol{u}}_{i}^{\mathrm{T}} & \dot{\boldsymbol{u}}_{j}^{\mathrm{T}} & \dot{\boldsymbol{u}}_{k}^{\mathrm{T}} \end{bmatrix} \frac{\rho h \Delta}{12} \begin{bmatrix} 2l_{2\times2} & l_{2\times2} & l_{2\times2} \\ l_{2\times2} & 2l_{2\times2} & l_{2\times2} \\ l_{2\times2} & l_{2\times2} & 2l_{2\times2} \end{bmatrix} \begin{bmatrix} \dot{\boldsymbol{u}}_{i} \\ \dot{\boldsymbol{u}}_{j} \\ \dot{\boldsymbol{u}}_{k} \end{bmatrix}$ 

(see Chapter\_2\_Finite\_Element\_Approximation.pdf at www.ritsumei.ac.jp/~hirai/edu/common/soft\_robotics/)

# Partial inertia matrix $M_{i,j,k} = \frac{\rho h \triangle}{12} \begin{bmatrix} 2l_{2\times2} & l_{2\times2} & l_{2\times2} \\ l_{2\times2} & 2l_{2\times2} & l_{2\times2} \\ l_{2\times2} & l_{2\times2} & 2l_{2\times2} \end{bmatrix}$



assume ho h riangle / 12 is constantly equal to 1 partial inertia matrices

$$M_{1,2,4} = M_{2,3,5} = M_{5,4,2} = M_{6,5,3} = \begin{bmatrix} 2I_{2\times2} & I_{2\times2} & I_{2\times2} \\ I_{2\times2} & 2I_{2\times2} & I_{2\times2} \\ I_{2\times2} & I_{2\times2} & 2I_{2\times2} \end{bmatrix}$$

#### total kinetic energy



*M*: inertia matrix ( $6 \times 6$  block matrix)

$$M_{1,2,4} = \begin{bmatrix} (1,1) \text{ block } (1,2) \text{ block } (1,4) \text{ block } \\ \hline (2,1) \text{ block } (2,2) \text{ block } (2,4) \text{ block } \\ \hline (4,1) \text{ block } (4,2) \text{ block } (4,4) \text{ block } \end{bmatrix}$$

contribution of  $M_{1,2,4}$  to M

Γ	$2I_{2\times 2}$	$I_{2\times 2}$	$I_{2\times 2}$	-
	$I_{2\times 2}$	$2I_{2\times 2}$	$I_{2\times 2}$	
-				
	$I_{2\times 2}$	$I_{2\times 2}$	$2I_{2\times 2}$	

$$M_{2,3,5} = \begin{bmatrix} (2,2) \text{ block} & (2,3) \text{ block} & (2,5) \text{ block} \\ \hline (3,2) \text{ block} & (3,3) \text{ block} & (3,5) \text{ block} \\ \hline (5,2) \text{ block} & (5,3) \text{ block} & (5,5) \text{ block} \end{bmatrix}$$

contribution of  $M_{2,3,5}$  to M

_				-
	$2I_{2\times 2}$	$I_{2\times 2}$	$I_{2\times 2}$	
	$I_{2\times 2}$	$2I_{2\times 2}$	$I_{2\times 2}$	
	$I_{2\times 2}$	$I_{2\times 2}$	$2I_{2\times 2}$	

$$M_{5,4,2} = \begin{bmatrix} (5,5) \text{ block} & (5,4) \text{ block} & (5,2) \text{ block} \\ \hline (4,5) \text{ block} & (4,4) \text{ block} & (4,2) \text{ block} \\ \hline (2,5) \text{ block} & (2,4) \text{ block} & (2,2) \text{ block} \end{bmatrix}$$

contribution of  $M_{5,4,2}$  to M

			-
$2I_{2\times 2}$	$I_{2\times 2}$	$I_{2\times 2}$	
$I_{2\times 2}$	$2I_{2\times 2}$	$I_{2\times 2}$	
 $I_{2\times 2}$	$I_{2\times 2}$	$2I_{2\times 2}$	

$$M_{6,5,3} = \begin{bmatrix} (6,6) \text{ block} & (6,5) \text{ block} & (6,3) \text{ block} \\ \hline (5,6) \text{ block} & (5,5) \text{ block} & (5,3) \text{ block} \\ \hline (3,6) \text{ block} & (3,5) \text{ block} & (3,3) \text{ block} \end{bmatrix}$$

#### contribution of $M_{6,5,3}$ to M

Γ				
		$2I_{2\times 2}$	$I_{2\times 2}$	$I_{2\times 2}$
		$I_{2\times 2}$	$2I_{2\times 2}$	$I_{2\times 2}$
		$I_{2\times 2}$	$I_{2\times 2}$	$2I_{2\times 2}$

inertia matrix

$$M = M_{1,2,4} \oplus M_{2,3,5} \oplus M_{5,4,2} \oplus M_{6,5,3}$$

$$= \begin{bmatrix} 2I_{2\times2} & I_{2\times2} & I_{2\times2} \\ I_{2\times2} & 6I_{2\times2} & I_{2\times2} & 2I_{2\times2} & I_{2\times2} \\ I_{2\times2} & 4I_{2\times2} & 2I_{2\times2} & I_{2\times2} \\ I_{2\times2} & 2I_{2\times2} & 4I_{2\times2} & I_{2\times2} \\ I_{2\times2} & 2I_{2\times2} & I_{2\times2} & I_{2\times2} \\ I_{2\times2} & 2I_{2\times2} & I_{2\times2} & I_{2\times2} \\ I_{2\times2} & I_{2\times2} & I_{2\times2} & 2I_{2\times2} \end{bmatrix}$$

assume  $\lambda$ ,  $\mu$  and h are constants strain potential energy stored in  $\triangle = \triangle P_i P_j P_k$ 

$$egin{aligned} & U_{i,j,k} = \int_{ riangle} rac{1}{2} \, oldsymbol{arepsilon}^{\mathrm{T}} (\lambda I_{\lambda} + \mu I_{\mu}) oldsymbol{arepsilon} \, h \, \mathrm{d}S \ &= rac{1}{2} \left[ egin{aligned} & oldsymbol{u}_i^{\mathrm{T}} & oldsymbol{u}_k^{\mathrm{T}} \end{array} 
ight] \, oldsymbol{\mathcal{K}}_{i,j,k} \, \left[ egin{aligned} & oldsymbol{u}_i \ & oldsymbol{u}_j \ & oldsymbol{u}_k \end{array} 
ight] \end{aligned}$$

where

$$K_{i,j,k} = \lambda J_{\lambda}^{i,j,k} + \mu J_{\mu}^{i,j,k}$$

(see Chapter\_2\_Finite\_Element\_Approximation.pdf)

$$oldsymbol{a} = rac{1}{2 riangle} \left[ egin{array}{c} y_j - y_k \ y_k - y_i \ y_i - y_j \end{array} 
ight], \qquad oldsymbol{b} = rac{-1}{2 riangle} \left[ egin{array}{c} x_j - x_k \ x_k - x_i \ x_i - x_j \end{array} 
ight]$$

$$egin{array}{ll} \mathcal{H}_{\lambda} = \left[egin{array}{ccc} m{aa}^{\mathrm{T}} & m{ab}^{\mathrm{T}} \ m{ba}^{\mathrm{T}} & m{bb}^{\mathrm{T}} \end{array}
ight] higtriangle \ \mathcal{H}_{\mu} = \left[egin{array}{ccc} 2m{aa}^{\mathrm{T}} + m{bb}^{\mathrm{T}} & m{ba}^{\mathrm{T}} \ m{ab}^{\mathrm{T}} & 2m{bb}^{\mathrm{T}} + m{aa}^{\mathrm{T}} \end{array}
ight] higtriangle \end{array}$$

1, 4, 2, 5, 3, 6 rows and columns of  $H_{\lambda}$ ,  $H_{\mu} \rightarrow$ 1, 2, 3, 4, 5, 6 rows and columns of  $J_{\lambda}^{i,j,k}$ ,  $J_{\mu}^{i,j,k}$  Example (stiffness matrix)



assume h = 2

stiffness matrix

$$K = K_{1,2,4} \oplus K_{2,3,5} \oplus K_{5,4,2} \oplus K_{6,5,3}$$
assume  $\lambda$  and  $\mu$  are constants over region

$$\begin{split} & \mathcal{K} = \mathcal{K}_{1,2,4} \oplus \mathcal{K}_{2,3,5} \oplus \mathcal{K}_{5,4,2} \oplus \mathcal{K}_{6,5,3} \\ &= (\lambda J_{\lambda}^{1,2,4} + \mu J_{\mu}^{1,2,4}) \oplus (\lambda J_{\lambda}^{2,3,5} + \mu J_{\mu}^{2,3,5}) \oplus \cdots \\ &= \lambda (J_{\lambda}^{1,2,4} \oplus J_{\lambda}^{2,3,5} \oplus \cdots) + \mu (J_{\mu}^{1,2,4} \oplus J_{\mu}^{2,3,5} \oplus \cdots) \\ &= \lambda J_{\lambda} + \mu J_{\mu} \end{split}$$

where

$$\begin{split} J_{\lambda} &= J_{\lambda}^{1,2,4} \oplus J_{\lambda}^{2,3,5} \oplus J_{\lambda}^{5,4,2} \oplus J_{\lambda}^{6,5,3} \\ J_{\mu} &= J_{\mu}^{1,2,4} \oplus J_{\mu}^{2,3,5} \oplus J_{\mu}^{5,4,2} \oplus J_{\mu}^{6,5,3} \end{split}$$





$$J_{\lambda}^{1,2,4} = J_{\lambda}^{2,3,5} = J_{\lambda}^{5,4,2} = J_{\lambda}^{6,5,3}$$

$$J^{1,2,4}_{\mu} = J^{2,3,5}_{\mu} = J^{5,4,2}_{\mu} = J^{6,5,3}_{\mu}$$

contribution of  $J_{\lambda}^{1,2,4}$  to  $J_{\lambda}$ 

Γ 1	1	-1	0	0	-1	-
1	1	-1	0	0	-1	
-1	-1	1	0	0	1	
0	0	0	0	0	0	
0	0	0	0	0	0	
-1	-1	1	0	0	1	
L						_

contribution of  $J_{\lambda}^{2,3,5}$  to  $J_{\lambda}$ 



contribution of  $J_{\lambda}^{5,4,2}$  to  $J_{\lambda}$ 



contribution of  $J_{\lambda}^{6,5,3}$  to  $J_{\lambda}$ 





$J_{\mu} = J_{\mu}^{1,2,4} \oplus J_{\mu}^{2,3,5} \oplus J_{\mu}^{5,4,2} \oplus J_{\mu}^{6,5,3}$												
=	- 3	1	-2	-1			-1	0				٦
	1	3	0	-1			-1	-2				
	-2	0	6	1	-2	-1	0	1	-2	-1		
	-1	-1	1	6	0	-1	1	0	-1	-4		
			-2	0	3	0			0	1	-1	-1
			-1	-1	0	3			1	0	0	-2
	-1	-1	0	1			3	0	-2	0		
	0	-2	1	0			0	3	-1	-1		
			-2	-1	0	1	-2	-1	6	1	-2	0
			-1	-4	1	0	0	-1	1	6	-1	-1
					-1	0			-2	-1	3	1
	_				-1	-2			0	-1	1	3 ]

stiffness matrix

$$\mathbf{K} = \lambda \mathbf{J}_{\lambda} + \mu \mathbf{J}_{\mu}$$

$$egin{array}{ccc} \lambda,\,\mu & {
m material-specific} \ J_\lambda,\,J_\mu & {
m geometric} \end{array}$$

strain potential energy

$$U=rac{1}{2}oldsymbol{u}_{\mathrm{N}}^{\mathrm{T}}$$
 K  $oldsymbol{u}_{\mathrm{N}}$ 

### Lagrange equation

Kinetic and strain potential energies

$$\mathcal{T} = rac{1}{2} \dot{\boldsymbol{\textit{u}}}_{\mathrm{N}}^{\mathrm{T}} \; \mathcal{M} \; \dot{\boldsymbol{\textit{u}}}_{\mathrm{N}}, \qquad \mathcal{U} = rac{1}{2} \boldsymbol{\textit{u}}_{\mathrm{N}}^{\mathrm{T}} \; \mathcal{K} \; \boldsymbol{\textit{u}}_{\mathrm{N}}$$

Work done by external forces

$$W = \boldsymbol{f}^{\mathrm{T}} \boldsymbol{u}_{\mathrm{N}}$$

Constraints

$$oldsymbol{R}=oldsymbol{A}^{\mathrm{T}}oldsymbol{u}_{\mathrm{N}}-oldsymbol{b}(t)=oldsymbol{0}$$

Lagrangian

$$\mathcal{L} = T - U + W + \lambda^{\mathrm{T}} R$$

# Lagrange equation Lagrange equation of motion and deformation $\frac{\partial \mathcal{L}}{\partial \boldsymbol{\mu}} - \frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial \mathcal{L}}{\partial \dot{\boldsymbol{\mu}}} = \boldsymbol{0}$ 11 $-Ku_{\rm N}+f+A\lambda-M\ddot{u}_{\rm N}=0$ $\downarrow$ $\dot{\boldsymbol{u}}_{\mathrm{N}} = \boldsymbol{v}_{\mathrm{N}}$

 $M\dot{v}_{
m N} - A\lambda = -Ku_{
m N} + f$ 

### Lagrange equation Canonical form

$$\dot{oldsymbol{u}}_{\mathrm{N}} = oldsymbol{v}_{\mathrm{N}} \ \begin{bmatrix} oldsymbol{M} & -A \ -A^{\mathrm{T}} \end{bmatrix} \begin{bmatrix} \dot{oldsymbol{v}}_{\mathrm{N}} \ oldsymbol{\lambda} \end{bmatrix} = \begin{bmatrix} -Koldsymbol{u}_{\mathrm{N}} + oldsymbol{f} \ oldsymbol{C}(oldsymbol{u}_{\mathrm{N}},oldsymbol{v}_{\mathrm{N}}) \end{bmatrix}$$

where

$$C(u_{\mathrm{N}}, v_{\mathrm{N}}) = -\ddot{b}(t) + 2\alpha (A^{\mathrm{T}}v_{\mathrm{N}} - \dot{b}(t)) + \alpha^{2} (A^{\mathrm{T}}u_{\mathrm{N}} - b(t))$$

any ODE solver can be applied to the canonical form

two-dimentional finite element calculation on MATLAB

http://www.ritsumei.ac.jp/~hirai/edu/common/ soft\_robotics/Physics\_Soft\_Bodies.html

Classes : NodalPoint, Triangle, Body

```
classdef NodalPoint
    properties
        Coordinates;
        Displacement;
        Velocity
    end
    methods
        function obj = NodalPoint(p)
            obj.Coordinates = p;
        end
    end
end
```

```
classdef Triangle
    properties
        Vertices;
        Area;
        Thickness:
        Density; lambda; mu;
        vector_a; vector_b;
        u_x; u_y; v_x; v_y;
        Cauchy_strain;
        Green_strain;
        Partial J lambda: Partial J mu:
        Partial_Stiffness_Matrix;
        Partial_Inertia_Matrix;
        Partial_Gravitational_Vector;
    end
   methods
```

```
classdef Body
      properties
           numNodalPoints; NodalPoints;
           numTriangles; Triangles;
           strain_potential_energy;
           gravitational_potential_energy;
           J_lambda; J_mu;
           Stiffness_Matrix;
           Inertia_Matrix;
           Gravitational_Vector;
      end
      methods
           function obj = Body(npoints, points, ntris, tris,
               obj.numNodalPoints = npoints;
               for k=1:npoints
                    pt(k) = NodalPoint(points(:,k));
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```

methods of class Triangle

partial\_derivaties calculating partial derivatives  $\partial u/\partial x$ ,  $\partial u/\partial y$ ,  $\partial v/\partial x$ ,  $\partial v/\partial y$ 

calculate\_Cauchy\_strain calculating Cauchy strain in the triangle partial\_strain\_potential\_energy strain potential energy stored in the triangle

calculate\_Green\_strain calculating Green strain in the triangle partial\_strain\_potential\_energy\_Green\_strain strain potential energy using Green strain

partial\_gravitational\_potential\_energy gravitational potential energy stored in the triangle

partial\_stiffness\_matrix calculating partial stiffness matrix  $K_{i,j,k}$ partial\_inertia\_matrix calculating partial inertia matrix  $M_{i,j,k}$ partial\_gravitational\_vector calculating partial gravitational vector

methods of class Body

total\_strain\_potential\_energy calculating strain potential energy stored in the body

- total\_strain\_potential\_energy\_Green\_strain strain potential energy using Green strain
- total\_gravitational\_potential\_energy gravitational potential energy stored in the body

calculate\_stiffness\_matrix calculating stiffness matrix Kcalculate\_inertia\_matrix calculating inertia matrix Mcalculate\_gravitational\_vector calculating gravitational vector gconstraint\_matrix constraint matrix when specified nodal points are fixed

draw draw the shape of the body

#### Example (dynamic simulation) two-dimensional square soft body of width wYoung's modulus E, viscous modulus c, density $\rho$ divide square into $3 \times 3 \times 2$ triangles



### Example (dynamic simulation) [0, $t_{push}$ ] fix the bottom & push P<sub>14</sub>P<sub>15</sub> downward [ $t_{push}$ , $t_{hold}$ ] fix the bottom & keep P<sub>14</sub>P<sub>15</sub> [ $t_{hold}$ , $t_{end}$ ] fix the bottom & release P<sub>14</sub>P<sub>15</sub>





Example (dynamic simulation)  $[0, t_{push}]$ note  $egin{array}{c|c} & oldsymbol{u}_1 & oldsymbol{u}_2 & oldsymbol{u}_2 & oldsymbol{u}_3 & oldsymbol{u}_4 & oldsymbol{u}_{14} & oldsymbol{u}_{14} & oldsymbol{u}_{15} & o$ 

speficies nodal points under constraints

# Example (dynamic simulation) [0, t<sub>push</sub>]

$$\boldsymbol{b}(t) = \boldsymbol{b}_0 + \boldsymbol{b}_1 t$$

where

$$\boldsymbol{b}_{0} = \begin{bmatrix} \boldsymbol{0} \\ \boldsymbol{0} \\ \boldsymbol{0} \\ \boldsymbol{0} \\ \boldsymbol{0} \\ \boldsymbol{0} \end{bmatrix}, \quad \boldsymbol{b}_{1} = \begin{bmatrix} \boldsymbol{0} \\ \boldsymbol{0} \\ \boldsymbol{0} \\ \boldsymbol{0} \\ \boldsymbol{v}_{push} \\ \boldsymbol{v}_{push} \end{bmatrix}$$
  
note  $\dot{\boldsymbol{b}}(t) = \boldsymbol{b}_{1}$  and  $\ddot{\boldsymbol{b}}(t) = \boldsymbol{0}$ , yielding  
 $\boldsymbol{C}(\boldsymbol{u}_{\mathrm{N}}, \boldsymbol{v}_{\mathrm{N}}) = 2\alpha(\boldsymbol{A}^{\mathrm{T}}\boldsymbol{v}_{\mathrm{N}} - \boldsymbol{b}_{1}) + \alpha^{2}(\boldsymbol{A}^{\mathrm{T}}\boldsymbol{u}_{\mathrm{N}} - (\boldsymbol{b}_{0} + \boldsymbol{b}_{1}t))$ 

### Example (dynamic simulation) [t<sub>push</sub>, t<sub>hold</sub>]



### Example (dynamic simulation) [t<sub>hold</sub>, t<sub>end</sub>]

$$oldsymbol{u}_1=oldsymbol{u}_2=oldsymbol{u}_3=oldsymbol{u}_4=oldsymbol{0}$$



% Dynamic deformation of an elastic square object (4&time % g, cm, msec

addpath('../two\_dim\_fea'); addpath('../two\_dim\_static'); width = 30; height = 30; thickness = 1; m = 4; n = 4; [points, triangles] = rectangular\_object(m, n, width, height)

```
Young = 10.0; c = 0.04*Young; nu = 0.48; density = 1.00;
Epfloor = 0.02; cpfloor = 0;
[lambda, mu] = Lame_constants(Young, nu);
[lambdav, muv] = Lame_constants(c, nu);
```

npoints = size(points,2);

```
% holding top region
b0 = [ zeros(2*4,1); 0; -vpush*tp; 0; -vpush*tp ];
b1 = zeros(2*6,1);
interval = [tp, tp+th];
qinit = q_push(end,:);
square_object_hold = @(t,q) square_object_constraint_para
[time_hold, q_hold] = ode23tb(square_object_hold, interval)
```

```
% releasing top region
A = elastic.constraint_matrix([1,2,3,4]);
b0 = zeros(2*4,1);
b1 = zeros(2*4,1);
interval = [tp+th, tp+th+tf];
qinit = q_hold(end,:);
square_object_free = @(t,q) square_object_constraint_para
[time_free, q_free] = ode23tb(square_object_free, interval)
```



#### simulation movie



#### simulation movie

#### Example (dynamic simulation) two-dimensional square soft body of width wYoung's modulus E, viscous modulus c, density $\rho$ divide square into $3 \times 3 \times 2$ triangles



[0,  $t_{push}$ ] fix the bottom & push  $P_{14}P_{15}$  downward [ $t_{push}$ ,  $t_{hold}$ ] fix the bottom & keep  $P_{14}P_{15}$ [ $t_{hold}$ ,  $t_{end}$ ] free (reaction force by penalty method)



% Jumping of an elastic square object (4×4) % g, cm, msec

addpath('../two\_dim\_fea'); addpath('../two\_dim\_static');

```
width = 30; height = 30; thickness = 1;
m = 4; n = 4;
[points, triangles] = rectangular_object(m, n, width, height)
```

```
Young = 10.0; c = 0.04*Young; nu = 0.48; density = 1.00;
Epfloor = 0.02; cpfloor = 0;
[lambda, mu] = Lame_constants(Young, nu);
[lambdav, muv] = Lame_constants(c, nu);
```

npoints = size(points,2);

```
% holding top region
b0 = [ zeros(2*4,1); 0; -vpush*tp; 0; -vpush*tp ];
b1 = zeros(2*6,1);
interval = [tp, tp+th];
qinit = q_push(end,:);
square_object_hold = @(t,q) square_object_constraint_para
[time_hold, q_hold] = ode23tb(square_object_hold, interval)
```

```
% releasing all constraints
floor_force = @(t,npoints,un,vn) floor_force_param(t,npoi
interval = [tp+th, tp+th+tf];
qinit = q_hold(end,:);
square_object_free = @(t,q) square_object_free_param(t,q,
[time_free, q_free] = ode23tb(square_object_free, interval)
```

```
time = [time_push; time_hold; time_free];
```
# Example (dynamic simulation)



jump simulation movie

# Example (dynamic simulation)



# Example (dynamic simulation)

- motion and deformation can be simulated properly
- results depend on mesh and include artifacts
- finer mesh yields better result but needs more computation time

### Handouts

Sample programs (MATLAB) are available at:

http://www.ritsumei.ac.jp/~hirai/edu/common/ soft\_robotics/Physics\_Soft\_Bodies.html

#### Report due date: 23:59, July 25 (Tues) Simulate the deformation of an elastic body $P_2P_3$ is fixed to the floor. $P_1$ and $P_4$ may slide on the floor. $[0, t_{push}]$ push $P_{14}P_{15}$ downward $[t_{push}, t_{hold}]$ keep $P_{14}P_{15}$ 13 14 15 16 $[t_{hold}, t_{end}]$ release $P_{14}P_{15}$ 10 9 12 11

