

# Sigma Models with Nontrivial Flux and Generalized Geometry

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## 1. Introduction

String theory (M theory) with

non zero fluxes

especially

non zero NS flux  $H \neq 0$

becomes more important

in the context of  $\rightarrow$

◦ flux compactification  $10D \rightarrow 4D$

NS flux  $H \neq 0$

RR fluxes  $F \neq 0$

vacua

provide new terms in

Kähler & super potentials  
in 4D theory

→ new phenomenological models

◦ direct generalization

of non commutative geometry

in string theory

NGG  $\left\{ \begin{array}{l} H = db = 0 \\ b \neq 0 \end{array} \right.$

NS B-field '99 Seiberg-Witten

↓  
GG  $H = db \neq 0$

In order to discuss  $H \neq 0$ ,

We need new geometry

$H=0$

$H \neq 0$

Complex

→

generalized complex

Kähler

→

generalized Kähler

Calabi-Yau

→

generalized Calabi-Yau

|

|

generalized geometry

'02 Hitchin

'04 Gualtieri

# Purpose of the project

is to investigate  
(mathematical / theoretical) aspects of

generalized geometry

in } string theory  
} M, F

→ Classification of string vacua  
Landscape

→ mathematical analysis

→ M theory

In this talk,

We are interested in

world sheet

(world volume)

descriptions

or

Sigma model

(membrane)

of generalized geometry

§ 2  $N=(2,2)$  SUSY sigma model

§ 3 definition of generalized complex structure

§ 4 topological sigma model  
topological membrane

• type II superstring 4 directions

non-topological

topological

worksheet 2D

topological twist

$N=(2,2)$  SUSY

$\sigma$ -model with

$H \neq 0$

'84 Gates, Hull, Roček

'04 Gualtieri

→ topological  $\sigma$ -model

With generalized CY

(generalized A-model

" B-model

'03 Kapustin, Li

'04 Zucchini

↑ '05 Pestun, Witten

target space 10D

$N=2$  Superstring

(SUGRA)

$H \neq 0$

'05 Grana, Minasian

Petrini, Tomasiello

et.al.

Review '06 Grana

6D

'03 Hitchin

Hitchin functional

(generalized)

With  $SU(3)$ -structure

'04 De Jonghe, Gaikov

Nestor, Vafa

'04 Nekrasov

'05 Pestun, Witten

• M-theory  $N=1$

non-topological

world volume 3D  
supermembrane

topological

3D  
topological 2-brane  
with  $G_2$  structure  
'05 Bonetti, Zerbini

2D  
 $N=1$   
String with  $G_2$  structure  
'91 Howe, Papadopoulos  
'94 Shatashvili, Vafa

? 2D  
topological  $G_2$  string  
'05 de Boer, Nagai, Skovran

target space

11D SUGRA

$G_{MN}$

$A_{MN}$  ↔ M2-brane  
M5-brane

7D  
Hitchin functional  
with  $G_2$  structure

'04 DGV

'04 Nekrasov

## 2. 2D $N=(2,2)$ SUSY Sigma Model with $H \neq 0$

'84 Gates, Hull, Roček

First we write manifest  $N=(1,1)$  SUSY  $\sigma$ -model

$$S = \frac{1}{2} \int_{\Sigma} d\sigma d\theta (g_{ij}(\Phi) + b_{ij}(\Phi)) D_{\pm} \Phi^i D_{\pm} \Phi^j$$

$$\bar{\Phi}^i \equiv \Phi^i + \theta^+ \psi_+^i + \theta^- \psi_-^i + \theta^- \theta^+ F^i$$

$N=(1,1)$  Superfield

$$\Phi^i: \Sigma \rightarrow M$$

2D world target sp.

$$H_{ijk} \equiv \frac{1}{3!} \partial_{[i} b_{jk]} \quad \text{NS flux}$$

$$g_{ij} = g_{ji} \quad \text{metric}$$

$$b_{ij} = -b_{ji} \quad \text{NS b-field}$$

$$D_{\pm} \equiv \frac{\partial}{\partial \theta^{\pm}} + i \theta^{\pm} \partial_{\pm} \quad \text{superderivative}$$



S has manifest  $N=(1,1)$  SUSY

$$\delta_1 \Phi^i \equiv \epsilon_1^+ D_+ \Phi^i + \epsilon_1^- D_- \Phi^i$$

If S has another SUSY

$$\delta_2 \Phi^i \equiv \epsilon_2^+ D_+ \Phi^i J_{+d}^i(\Phi) + \epsilon_2^- D_- \Phi^i J_{-d}^i(\Phi)$$

$(g, b, J_{\pm})$  must have some conditions  
then S has  $N=(2,2)$  SUSY.

② Conditions

• If  $H=0$

M is Kähler  $\Rightarrow N=(2,2)$  SUSY

$J_{\pm}$  : complex structure

$\omega_{\pm} \equiv g J_{\pm}$  : Kähler form

• If  $H \neq 0$

①  $J_{\pm}$  are integrable almost complex str.

②  $g_{ij}$  is hermitian w.r.t  $J_+$  &  $J_-$

$$\textcircled{3} \quad \nabla_i^{(\pm)} J_{\pm}^j \equiv \partial_i J_{\pm}^j - T_{\pm}^j{}_{\ell} J_{\pm}^{\ell} - T_{\pm}^{\ell}{}_{i} J_{\pm}^j = 0$$

where

$$T_{\pm}^i{}_{jk} \equiv T_{\pm}^i{}_{jk} \pm \frac{1}{2} g^{il} H_{\ell jk}$$

covariant derivative with torsion

①②③ is called  
the bi-Hermitian structure.

Generally  $\omega_{\pm} \equiv g J_{\pm}$  is not closed.  
(non Kähler)

But this formulation is not so useful  
because the relation to the Kähler  
geometry is not manifest.

bi-Hermitian structure ← '04 Gualtieri  
 $\Leftrightarrow$  (twisted) generalized Kähler str.

### 3. Generalized Geometry (generalized complex structure)

cf. We remember the definition of the complex structure on  $M$

$M$ :  $d$ -dim manifold

$TM$ : tangent bundle

$$J: TM \rightarrow TM \quad \text{map}$$

$$\left( J_{\delta}^i: X^i \rightarrow J_{\delta}^i X^{\delta} \right)$$

S.t.

$$J^2 = -1 \quad \text{--- (1)}$$

$\pm i$  are eigenvalues of  $J$ .

Let  $\Pi_{\pm}$  be projections on  $\pm i$  eigen bundle.

Integrable condition is

$$\Pi_{\mp} [\Pi_{\pm} X, \Pi_{\pm} Y] = 0 \quad \text{--- (2)}$$

where  $[\ , \ ]$ : Lie bracket  
 $X, Y \in TM$

## ⊙ generalized complex structure

We replace  $TM$  to  $TM \oplus T^*M$

$X, Y, \dots$  vector field on  $TM$

$\xi, \eta, \dots$  1-form on  $T^*M$

$X + \xi$  formal sum  $\in TM \oplus T^*M$

⊙ Inner product  $\langle \cdot, \cdot \rangle$

$$\langle X + \xi, Y + \eta \rangle \equiv \frac{1}{2} (i_X \eta + i_Y \xi)$$

$i_X, i_Y$ : interior product

$$\text{i.e. } \left\langle X \frac{\partial}{\partial \phi^i} + \xi_i d\phi^i, Y \frac{\partial}{\partial \phi^j} + \eta_j d\phi^j \right\rangle = \frac{1}{2} (X \eta_j + \xi_i Y^i)$$

$\{\phi^i\}$  local coordinate on  $M$

⊙ signature  $(d, d)$

⊙ on orthogonal coordinate

$$I = \begin{pmatrix} 0 & 1_d \\ 1_d & 0 \end{pmatrix}$$

◦ Courant bracket

$$[X+\xi, Y+\eta] \equiv [X, Y] + \mathcal{L}_X \eta - \mathcal{L}_Y \xi - \frac{1}{2} d_M(\iota_X \eta - \iota_Y \xi)$$

$\mathcal{L}_X, \mathcal{L}_Y$  : Lie differential

$d_M$  : exterior derivative on  $M$

◦ antisymmetric

◦ does not satisfy the Jacobi identity

Note)  $\Leftrightarrow$  Dorfman bracket

$$(X+\xi) \circ (Y+\eta) \equiv [X, Y] + \mathcal{L}_X \eta - \iota_Y d\xi$$

◦ not antisymmetric

◦ satisfies the Jacobi identity

$$[X+\xi, Y+\eta] = (X+\xi) \circ (Y+\eta) - (Y+\eta) \circ (X+\xi)$$

Def generalized complex structure  $f$

$$f: TM \oplus T^*M \rightarrow TM \oplus T^*M$$

s.t.  $f$  preserves  $\langle, \rangle \Leftrightarrow f^* I f = I$  — ①

$$f^2 = -1$$

— ①

$$\Pi_{\mp} [\Pi_{\pm}(X+\xi), \Pi_{\pm}(Y+\eta)] = 0 \quad \text{--- ②}$$

◦  $\Pi_{\pm} \equiv \frac{1}{2} (1 \mp \sqrt{-1} f)$

projections on  $\pm \sqrt{-1}$  eigen bundles

◦  $[ , ]$  Courant bracket

We represent  $2 \times 2$  matrix

$$f = \begin{pmatrix} J_{ij} & P_{ij} \\ Q_{ij} & -J_{ij}^* \end{pmatrix} \begin{matrix} TM \\ T^*M \\ TM \\ T^*M \end{matrix}$$

### local coordinate expression

$$f = \begin{pmatrix} J & P \\ Q & -J^* \end{pmatrix} \begin{matrix} TM \\ T^*M \end{matrix}$$

①

$$J^i_k J^k_j + P^{ik} Q_{kj} + \delta^i_j = 0$$

$$J^i_k P^{kj} + J^j_k P^{ki} = 0$$

$$Q_{ik} J^k_j + Q_{jk} J^k_i = 0$$

$$P^{ij} + P^{ji} = 0, Q_{ij} + Q_{ji} = 0$$

②

$$A^{ijk} = \beta_i^{jk} = C_{ij}^k = D_{ijk} = 0$$

$$A^{ijk} = P^{il} \partial_l P^{jk} + (ijk \text{ cyclic})$$

$$\beta_i^{jk} = J^l_i \partial_l P^{jk} + P^{il} (\partial_i J^k_l - \partial_l J^k_i) + P^{kl} \partial_l J^j_i - \partial_i J^j_l P^{lk}$$

$$C_{ij}^k = J^l_i \partial_l J^k_j - J^l_j \partial_l J^k_i - J^k_l \partial_i J^l_j + J^k_l \partial_j J^l_i + P^{kl} (\partial_l Q_{ij} + \partial_i Q_{jl} + \partial_j Q_{li})$$

$$D_{ijk} = J^l_i (\partial_l Q_{jk} + \partial_k Q_{lj}) + J^l_i (\partial_l Q_{ki} + \partial_i Q_{lk}) + J^l_k (\partial_l Q_{ij} + \partial_j Q_{li}) - Q_{jl} \partial_i J^l_k - Q_{ke} \partial_j J^l_i - Q_{ie} \partial_k J^l_j$$

example

$$1. J = \begin{pmatrix} J & 0 \\ 0 & -J^* \end{pmatrix}$$

$J$  is GCS  $\Leftrightarrow J$  is complex str.

$$2. J = \begin{pmatrix} 0 & -\Omega^{-1} \\ \Omega & 0 \end{pmatrix}$$

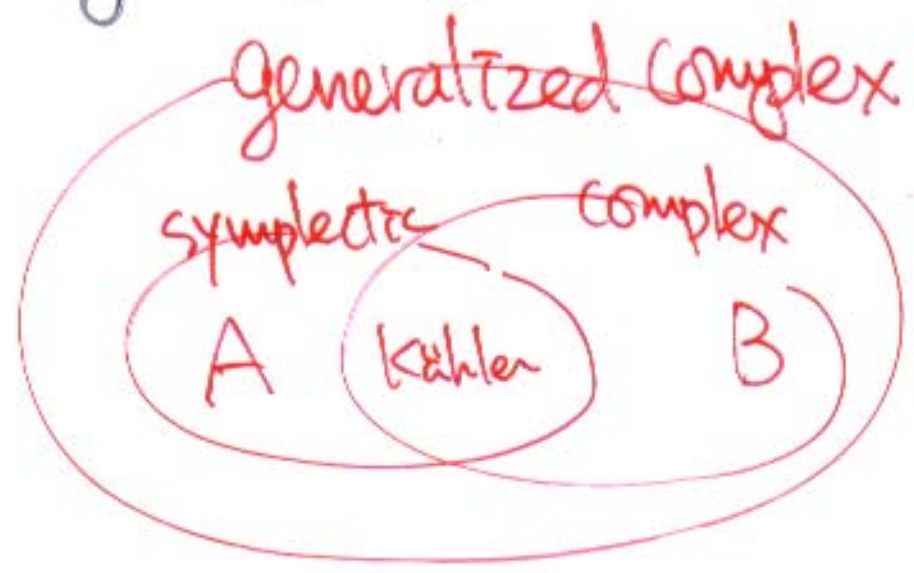
$J$  is GCS  $\Leftrightarrow \Omega$  is symplectic str.

### $\Omega$ twisted Generalized Complex structure

We can generalize GCS by a closed 3-form  $H$

$$[X+\zeta, Y+\eta]_H \equiv [X+\zeta, Y+\eta] + i_X i_Y H$$

In integrability condition (2),  $[\cdot, \cdot]$  is replaced by  $[\cdot, \cdot]_H$





then

② is modified

$$A_H{}^{\dot{i}\dot{j}k} = A^{\dot{i}\dot{j}k}$$

$$B_H{}^{\dot{i}k} = B_i{}^{\dot{i}k} + P^{il} P^{km} H_{lkm}$$

$$C_H{}^{\dot{i}k} = C_{\dot{i}j}{}^k - J^l{}_i P^{km} H_{jlm} \\ + J^l{}_j P^{km} H_{ilm}$$

$$D_H{}^{\dot{i}jk} = D_{\dot{i}jk} - H_{ijk} + J^l{}_i J^m{}_j H_{klm} \\ + J^l{}_j J^m{}_k H_{ilm} + J^l{}_k J^m{}_i H_{jlm}$$

$$A_H = B_H = C_H = D_H = 0$$

# Auto morphism on GCS

semi-direct product of

$Diff(M)$  &  $b$ -transformation

( $b$ -field transf.)

Def  $b$ -transformation  $(Diff(M) \times \Omega^2_{closed}(M))$

for 2-form  $b = b_{ij} d\phi^i d\phi^j$

$$\exp(b)(X+\xi) \equiv X+\xi + i_X b$$

If  $dm b = 0 \Rightarrow$

$$[\exp(b)(X+\xi), \exp(b)(Y+\eta)]$$

$$= \exp(b)[X+\xi, Y+\eta] \quad \text{covariant}$$

$$2. \hat{g} \equiv \exp(-b) g \exp(b) \quad \text{adjoint}$$

local coordinate expression

$$\hat{J}^i_j = J^i_j - P^{ik} b_{kj}$$

$$\hat{P}^{ij} = P^{ij}$$

$$\hat{Q}_{ij} = Q_{ij} + b_{ik} J^k_j - b_{jk} J^k_i + P^{kl} b_{ki} b_{lj}$$

$b$ -transformation

$$\begin{pmatrix} 1 & 0 \\ b & 1 \end{pmatrix}$$

• If  $H = dmb \neq 0$ .

$$\begin{aligned} & [\exp(b)(X+\xi), \exp(b)(Y+\eta)]_{H+dm} \\ &= \exp(b) [X+\xi, Y+\eta]_H \end{aligned}$$

$$H \xrightarrow{b\text{-transf}} \hat{H} = H - dmb$$

+ GCS is defined for  $H \in H^3(M)$

## Q generalized Kähler structure

$J_1, J_2$ : two <sup>(twisted)</sup> GCS

generalized Kähler

$\Leftrightarrow G \equiv -J_1 J_2$  is a positive definite metric on  $M$ .

example) Kähler manifold  $(g, J, \omega)$

$$J_1 = \begin{pmatrix} J & 0 \\ 0 & -J^* \end{pmatrix}, J_2 = \begin{pmatrix} 0 & \omega^{-1} \\ \omega & 0 \end{pmatrix}$$

then  $G = \begin{pmatrix} 0 & g^{-1} \\ g & 0 \end{pmatrix}$       $g = \omega J$

is metric.

---

We return to  $N=(2,2)$   $\sigma$ -model

$(g, b, J_{\pm})$ : bi-Hermitian str.

If we define

$$J_{1/2} \equiv \frac{1}{2} \begin{pmatrix} 1 & 0 \\ b & 1 \end{pmatrix} \begin{pmatrix} J_+ \pm J_- & -(\omega_+^{-1} \mp \omega_-^{-1}) \\ \omega_+ \mp \omega_- & -({}^t J_+ \pm {}^t J_-) \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -b & 1 \end{pmatrix}$$

then  $\omega_{\pm} \equiv g J_{\pm}$

$(g, b, J_{\pm})$  bi-Hermitian

$\Leftrightarrow J_1, J_2$  generalized Kähler

# 4. topological sigma model

We can construct a topological  $\sigma$ -model corresponding to  $N=(2,2)$   $\sigma$ -model in order to analyze only a vacuum (geometry).

⊗ topological twist <sup>1911</sup> <sup>'03 Kapustin</sup> <sup>written '04 Kapustin</sup> <sup>LI</sup>  
changes in  $N=(2,2)$   $\sigma$ -model

- $P_1, P_2$  translation
- $M$  Lorentz
- $Q_{\pm}, \bar{Q}_{\pm}$  SUSY
- $I_V$  vector R-sym
- $I_A$  axial R-sym

we replace Lorentz generator  $M$  to

$$M \rightarrow M' = M + \frac{1}{2} I_V \quad (\text{generalized}) \text{ A-twist}$$

then BRST  $Q_B = \frac{Q_+ + \bar{A}_+}{2} + \frac{Q_- - \bar{Q}_-}{2}$

$$M \rightarrow M' = M + \frac{1}{2} I_A \quad (\text{generalized}) \text{ B-twist}$$

BRST  $Q_B = \frac{Q_+ + \bar{Q}_+}{2} + \frac{Q_- + \bar{Q}_-}{2}$

In order to make topological twist  
 $I_V$  or  $I_A$  cannot have any anomaly

$\pm \sqrt{1}$  eigenbundles w.r.t  $J_+$   
 $T_+^{1,0} \oplus T_+^{0,1}$

$\pm \sqrt{-1}$  eigenbundles w.r.t  $J_-$   
 $T_-^{1,0} \oplus T_-^{0,1}$

Anomaly cancellation condition

$$U(1)_V : C_1(T_-^{1,0}) - C_1(T_+^{1,0}) = 0$$

$$U(1)_A : C_1(T_-^{1,0}) + C_1(T_+^{1,0}) = 0$$

$\Leftrightarrow$  **generalized Calabi-Yau condition**

where  $C_1(T)$  : 1-st Chern class

is a generalization of CY condition

Note)  $b=0$

If  $J_+ = J_-$  & A-twist  $\rightarrow$  A-model

B-twist  $\rightarrow$  B-model

'91 Witten

If  $b=0$

A-twist (A-model) **Kähler**

• need not CY condition

• depends only on  $\mathcal{J} = \begin{pmatrix} 0 & \omega' \\ \omega & 0 \end{pmatrix}$

B-twist (B-model) **Complex**

• does not have quantum correction

• depends only on  $\mathcal{J} = \begin{pmatrix} J & 0 \\ 0 & -J^* \end{pmatrix}$

**But**

If  $H=db \neq 0$

generalized A & B twist

do not have special properties

• depends on  $\mathcal{J} = \begin{pmatrix} J & P \\ Q & -J^* \end{pmatrix}$

Let us construct the topological twisted action in the case of  $H \neq 0$

Procedure

- Change fermion statistics  
(A or B-twist)
- ↓
- Rewrite the action

$$S = S_{\text{top}} + \delta_B(\ast)$$

by decoupling the BRST exact terms

'97 Alexandrov, Kontsevich, Schwarz, Zabzine  
'06 Bredthauer, Lindström, Persson, Zabzine

The answer is the Zucchini action

'04 Zucchini

$$\begin{aligned}
 S_{\text{top}} &= \int \left( B_i d\phi^i + J^i_j(\phi) B_i d\phi^j \right. \\
 &\quad \left. + \frac{1}{2} P^{ij}(\phi) B_i B_j + \frac{1}{2} Q_{ij}(\phi) d\phi^i d\phi^j \right) \\
 &= \int B_i d\phi^i + \frac{1}{2} \int (d\phi^i B_i) \int \begin{pmatrix} d\phi^j \\ B_j \end{pmatrix}
 \end{aligned}$$

$\phi$ : scalar superfield

$B_i \sim (g_{ij} + b_{ij}) d\phi^j$ : auxiliary 1-form superfield



However

$J, P, Q$  is GCS  $\Rightarrow$   $S_{\text{top}}$  is BRST inv.  
 $\Leftarrow$

Q If  $H \neq 0$  twisted GCS

We add the WZ term

$X$ : 3D mfd  
s.t.  $\partial X = \Sigma$



$$S'_{\text{top}} = S_{\text{top}} + \int_X H_{ijk} d\phi^i d\phi^j d\phi^k$$

$J, P, Q, H$  is twisted GCS

$\Rightarrow$   $S'_{\text{top}}$  is BRST inv.  
 $\Leftarrow$

Moreover

b-transformation

$$\begin{cases} \hat{\phi}^i = \phi^i \\ \hat{B}_i = B_i + b_{ij} d\phi^i d\phi^j \end{cases}$$

$$\hat{S}_{\text{top}} = S_{\text{top}} - \int_{\Sigma} b_{ij} d\phi^i d\phi^j \quad \text{not b-inv.}$$

## Q Our Purpose & Program

- Stop is not  $b$ -transformation inv.  
→  $b$ -transf. invariant formulation
- topological string .. String  
→ topological M theory 2-brane
- 'Courant bracket' is not familiar to physicists. →  
The algebra closed by the Courant bracket is the Courant algebroid.

{ Lie bracket → Lie algebra → gauge theory  
Courant bracket → Courant algebroid →

◦ 2D  
WZW model | Chern-Simons theory  
Zucchini Stop |

# Courant algebroid

vector bundle  $E \rightarrow M$

'90 Courant ~~132~~  
'97 Liu, Weinstein  
Xu

with  $\langle \cdot, \cdot \rangle$ : (graded) symmetric bilinear form

$\circ$ : bilinear form (Dorfman bracket)

$\rho: E \rightarrow TM$  anchor

s.t.

$$1. e_1 \circ (e_2 \circ e_3) = (e_1 \circ e_2) \circ e_3 + e_2 \circ (e_1 \circ e_3)$$

$$2. \rho(e_1 \circ e_2) = [\rho(e_1), \rho(e_2)]$$

$$3. e_1 \circ Fe_2 = F(e_1 \circ e_2) + (\rho(e_1)F)e_2$$

$$4. e_1 \circ e_2 = \frac{1}{2} \mathcal{D} \langle e_1, e_2 \rangle$$

$$5. \rho(e_1) \langle e_2, e_3 \rangle = \langle e_1 \circ e_2, e_3 \rangle + \langle e_2, e_1 \circ e_3 \rangle$$

where  $e_1, e_2, e_3 \in \Gamma(E)$

$F \in C^\infty(M)$

$\mathcal{D}: M \rightarrow \Gamma(E)$  s.t.  $\langle \mathcal{D}F, e \rangle = \rho(e)F$

We consider topological 2-brane  
 $X \rightarrow M$   
4-superfields

$\phi^i : X \rightarrow M$  scalar superfield

$A^i \in T^*X \otimes \phi^*(TM)$  1-form superfield

$B_{1i} \in T^*X \otimes \phi^*(T^*M)$  1-form "

$B_{2i} \in \wedge^2 T^*X \otimes \phi^*(T^*M)$  2-form "

We consider BF type kinetic term

$$S_0 = \int_X B_{2i} d\phi^i + B_{1i} dA^i$$

Most general interaction term  $S_1$

$$S = S_0 + g S_1$$

is

# Courant Sigma Model (30)

'02 N.I.

$$S = S_0 + g S_1$$

answer

$$S_1 = \int_X (f_1 \overset{i}{\phi} A_1^a B_{2i} + f_2 \overset{i}{\phi} B_{2i} B_{1b} + \frac{1}{3!} f_{3abc} \overset{i}{\phi} A_1^a A_1^b A_1^c + \frac{1}{2} f_{4ab} \overset{i}{\phi} A_1^a A_1^b B_{1c} + \frac{1}{2} f_{5a} \overset{i}{\phi} A_1^a B_{1b} B_{1c} + \frac{1}{3!} f_6 \overset{i}{\phi} B_{1a} B_{1b} B_{1c})$$

where  $f_1 \sim f_6$  satisfy

$$(S_1, S_1) = 0$$

↓ obstruction

$$\textcircled{1} f_{1e} \overset{i}{\phi} f_2^{\dot{e}} + f_2^{\dot{e}} f_{1e} \overset{i}{\phi} = 0$$

$$\textcircled{2} - \frac{\partial f_{1c} \overset{i}{\phi}}{\partial \phi^{\dot{a}}} f_{1b}^{\dot{a}} + \frac{\partial f_{1b} \overset{i}{\phi}}{\partial \phi^{\dot{a}}} f_{1c}^{\dot{a}} + f_{1e} \overset{i}{\phi} f_{4bc}^e + f_2^{\dot{e}} f_{3abc}^e = 0$$

$$\textcircled{3} f_{1b}^{\dot{a}} \frac{\partial f_2^{\dot{c}}}{\partial \phi^{\dot{a}}} - f_2^{\dot{c}} \frac{\partial f_{1b} \overset{i}{\phi}}{\partial \phi^{\dot{a}}} + f_{1e} \overset{i}{\phi} f_{5b}^e - f_2^{\dot{e}} f_{4eb}^c = 0$$

$$\textcircled{4} - f_2^{\dot{b}} \frac{\partial f_2^{\dot{c}}}{\partial \phi^{\dot{a}}} + f_2^{\dot{c}} \frac{\partial f_2^{\dot{b}}}{\partial \phi^{\dot{a}}} + f_{1e} \overset{i}{\phi} f_6^{\dot{e}bc} + f_2^{\dot{e}} f_{5e}^{\dot{b}c} = 0$$

$$\textcircled{5} - f_{1a} \overset{i}{\phi} \frac{\partial f_{4bc}^d}{\partial \phi^{\dot{a}}} + f_2^{\dot{d}} \frac{\partial f_{3abc}^d}{\partial \phi^{\dot{a}}} + f_{4e} \overset{d}{\phi} f_{abc}^e + f_3^{\dot{d}e} f_{5c}^{\dot{d}e} = 0$$

$$\textcircled{6} -f_{1a}^i \frac{\partial f_{5b}^{cd}}{\partial \phi^i} + f_2^i [c \frac{\partial f_{4ab}^{d]}{\partial \phi^i} + f_{3eab} f_6^{ecd} + f_{4e} [a f_{5b}^{c]e} + f_{4ab}^e f_{5e}^{cd}] = 0$$

$$\textcircled{7} -f_{1a}^i \frac{\partial f_6^{bcd}}{\partial \phi^i} + f_2^i [b \frac{\partial f_{5a}^{cd}}{\partial \phi^i} + f_{4ea} [b f_6^{cd}]e + f_{5e} [bc f_{5a}^{d}]e] = 0$$

$$\textcircled{8} -f_2^i [a \frac{\partial f_6^{bcd}}{\partial \phi^i} + f_6^e [ab f_{5e}^{cd}] = 0$$

$$\textcircled{9} -f_{1a}^i \frac{\partial f_{3bcd}}{\partial \phi^i} + f_{4[ab}^e f_{3cd]e} = 0$$

$\Leftrightarrow Q_B^2 = 0 \Leftrightarrow$  gauge symmetry  
 $\Leftrightarrow$  the Courant algebroid.

Courant algebroid  $\otimes \mathbb{C} \xrightarrow{\mathcal{J}} \text{GCS}$   
 $\mathcal{J} = \text{IFI projection}$

therefore

3D Schwarz type TFT (Courant  $\sigma$ -model)  
2-brane

$\downarrow \mathcal{J} = \text{IFI projection}$

3D generalized complex  $\sigma$ -model

Generalized Complex Sigma model in 3D

$\phi^i : X \rightarrow M$        $X : 3D$

$A^i \in T^*X \otimes \phi^*(T^*M)$       1-form

$B_{ij} \in T^*X \otimes \phi^*(T^*M)$       1-form

$B_{2i} \in \Lambda^2 T^*X \otimes \phi^*(T^*M)$       2-form

$$S_{GC} = \int_X -B_{2i} d\phi^i + B_{ij} dA^i - J^i_j B_{2i} A^j - P^{ij} B_{2i} B_{1j} + \frac{1}{2} (J^l_i H_{jke} + \frac{\partial Q_{jk}}{\partial \phi^i}) A^i A^j A^k + \frac{1}{2} (-P^{kl} H_{jle} - \frac{\partial J^k_j}{\partial \phi^i} + \frac{\partial J^k_i}{\partial \phi^j}) A^i A^j B_{1k} + \frac{1}{2} \frac{\partial P^{jk}}{\partial \phi^i} A^i B_{1j} B_{1k}$$

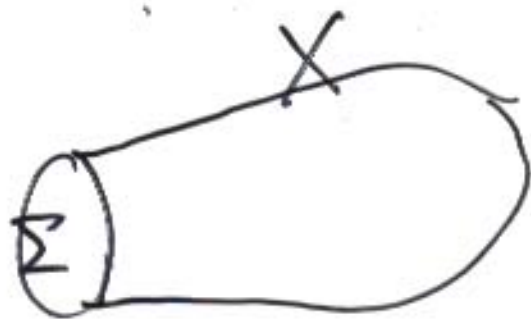
- $S_{GC}$  is BRST inv.  $\Leftrightarrow J, P, Q, H$  is (twisted) GCS
- $S_{GC}$  is b-transformation invariant.



## 6. Zucchini Model as a boundary action (16)

$$\Sigma = \partial X$$

$$\Sigma: 2D, X: 3D$$



We can derive <sup>2D</sup> Zucchini model as a boundary action on  $\Sigma$  of 3D GC sigma model.

H=0 case

$$S_{GC}(X) = S_{\Sigma}(\Sigma) + S_a(X)$$

- $S_{\Sigma}(\Sigma)$ : Zucchini action  $B_{ii} = B_i$
- $S_a(X)$ : topological term independent of GCS

$$S_a(X) = \int_X -Y_{2i}' d\phi^i + dB_{ii} Z'^i + Y_{2i} A^i + B_{2i} Z'^i + B_{2i} A^i$$

$Y_{2i}', Z'^i$ : auxiliary fields

$S_{GC}(X)$  &  $S_Z(\Sigma)$  is based on the same GCS, so  $S_Z(\Sigma)$  is regarded as boundary action of  $S_{GC}(X)$

$H \neq 0$  case

We can obtain

$$S'_{GC}(X) = S'_Z(\Sigma) + S_a(X)$$

$$S'_Z(\Sigma) = S_Z(\Sigma) + \frac{1}{2} \int_{\Sigma} T^l_{ij} H_{jka} d\phi^i d\phi^j d\phi^k - p^{k\ell} H_{ij\ell} B_{ik} d\phi^i d\phi^j$$

b-inv. WZ term

Zucchin action is not b-invariant

$$S_{GC} = S_Z(\Sigma) + S_a(X)$$

b-transformation

$$\hat{S}_{GC} = \hat{S}_Z(\Sigma) + \hat{S}_a(X)$$

$$\hat{S}_Z(\Sigma) = S_Z(\Sigma) - \frac{1}{2} \int_{\Sigma} b_{ij} d\phi^i d\phi^j$$

$H = -dmb$

$$= S_Z(\Sigma) + \frac{1}{2} \int_X H = \text{Zucchin model with } H$$

# Summary

3D topological  $\sigma$ -model  
with twisted GCS on  $X$

↓  $\partial X$  boundary

2D topological  $\sigma$ -model  
with twisted GCS on  $\partial X$

- BRST  $\Leftrightarrow$  tGCS
- b-transformation invariant

Future direction  $H \neq 0$

nontopological

topological

worldsheet

~~topological twist~~  
quantization  
CFT?



correspondence

target sp

M theory

relation?

super membrane



topological membrane